

Photon-induced processes in Pythia 8

ATLAS WORKSHOP FOR PHOTON-INDUCED PROCESSES

Ilkka Helenius

November 3rd, 2022



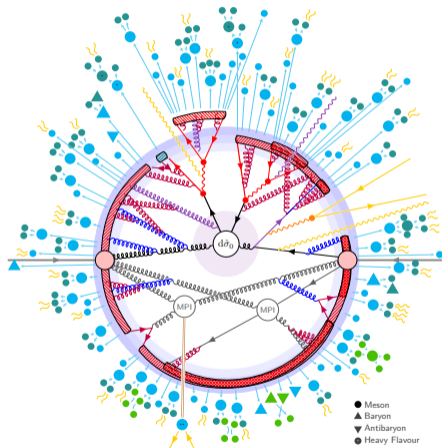
Outline

Pythia 8: A general purpose Monte-Carlo event generator

- A new manual for 8.3 release
arXiv:2203.11601 [hep-ph]

Outline

1. Photoproduction
2. Photon-induced processes in p+p
3. Ultraperipheral heavy-ion collisions
4. Diffractive photoproduction
5. Summary & Outlook



[figure by P. Skands]

Photoproduction

Events in e+p classified in terms of photon virtuality Q^2

- Large Q^2 : Deep inelastic scattering
- Small Q^2 : Photoproduction \Rightarrow Factorize photon flux

Direct processes

- Convolute photon flux f_γ with proton PDFs f_i^p and $d\hat{\sigma}$

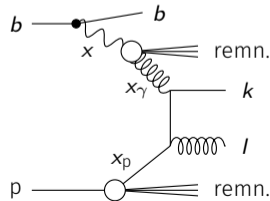
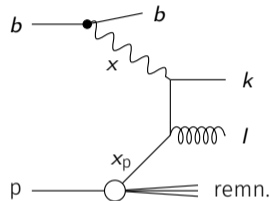
$$d\sigma^{bp \rightarrow kl+X} = f_\gamma^b(x) \otimes f_i^p(x_p, \mu^2) \otimes d\hat{\sigma}^{\gamma i \rightarrow kl}$$

Resolved processes

- Can fluctuate into a hadronic state: photon PDFs

$$d\sigma^{bp \rightarrow kl+X} = f_\gamma^b(x) \otimes f_j^\gamma(x_\gamma, \mu^2) \otimes f_i^p(x_p, \mu^2) \otimes d\sigma^{ij \rightarrow kl}$$

- Evolve γp as any hadronic collision (including MPIs)



Comparison to HERA dijet photoproduction data

ZEUS dijet measurement

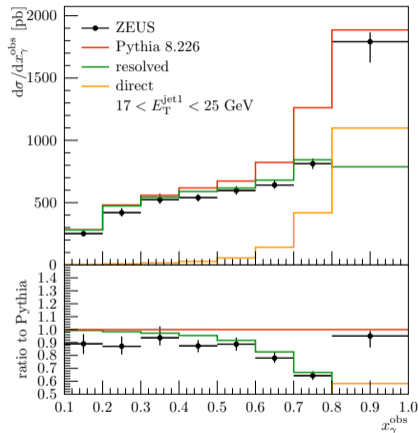
- $Q^2 < 1.0 \text{ GeV}^2$
- $134 < W_{\gamma p} < 277 \text{ GeV}$
- $E_T^{\text{jet1}} > 14 \text{ GeV}, E_T^{\text{jet2}} > 11 \text{ GeV}$
- $-1 < \eta^{\text{jet1,2}} < 2.4$

Two contributions

- Momentum fraction of partons in photon

$$x_\gamma^{\text{obs}} = \frac{E_T^{\text{jet1}} e^{\eta^{\text{jet1}}} + E_T^{\text{jet2}} e^{\eta^{\text{jet2}}}}{2yE_e} \approx x_\gamma$$

- Sensitivity to process type
- At high- x_γ^{obs} direct processes dominate



[ZEUS: Eur.Phys.J. C23 (2002) 615-631]

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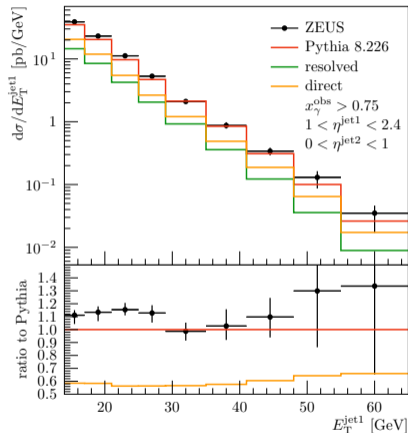
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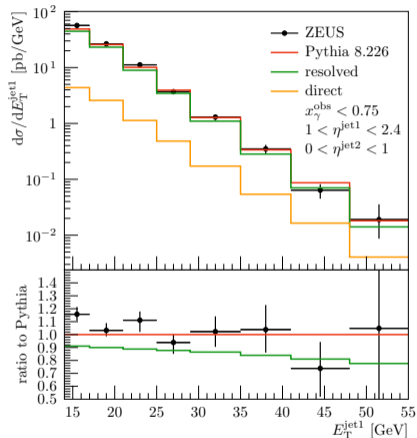
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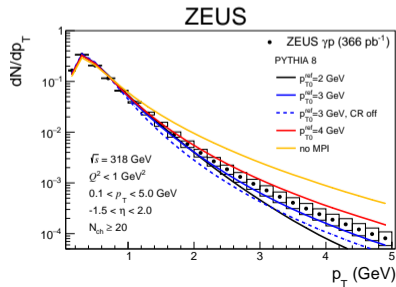
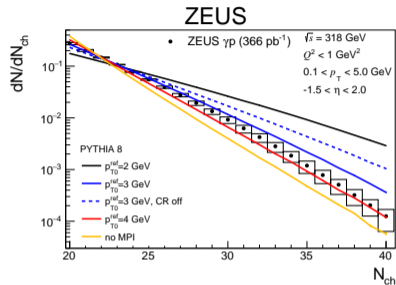
Multiplicity distributions

- Multiplicity distributions sensitive to MPIs with resolved photons
- ZEUS data support for MPIs but with slightly larger p_{T0}^{ref} than in pp \Rightarrow less MPIs

p_T spectra for $N_{\text{ch}} > 20$

- Similar agreement as above
- Useful constraints for MPIs in γp system
- Good agreement also in $c_1\{2\}$

[Rivet Analysis in preparation]



Equivalent photon approximation

Implemented photon fluxes

- In case of a point-like lepton we have

$$f_{\gamma}^l(x, Q^2) = \frac{\alpha_{em}}{2\pi} \frac{(1 + (1-x)^2)}{x} \frac{1}{Q^2}$$

- For protons need to account the form factor

$$f_{\gamma}^p(x, Q^2) = \frac{\alpha_{em}}{2\pi} \frac{(1 + (1-x)^2)}{x} \frac{1}{Q^2} \frac{1}{(1 + Q^2/Q_0^2)^4}$$

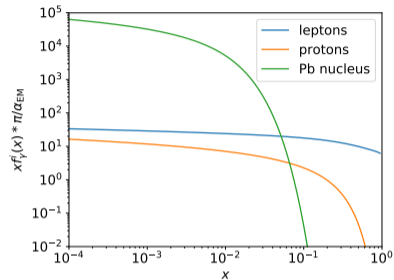
where $Q_0^2 = 0.71 \text{ GeV}^2$ (Drees-Zeppenfeld) \Rightarrow Large Q^2 heavily suppressed

- With heavy nuclei use b -integrated point-like-charge flux

$$f_{\gamma}^A(x) = \frac{2\alpha_{EM}Z^2}{x\pi} \left[\xi K_1(\xi)K_0(\xi) - \frac{\xi^2}{2} (K_1^2(\xi) - K_0^2(\xi)) \right]$$

where $\xi = b_{\min} x m$ where b_{\min} reject nuclear overlap, $Q^2 \ll 1 \text{ GeV}^2$

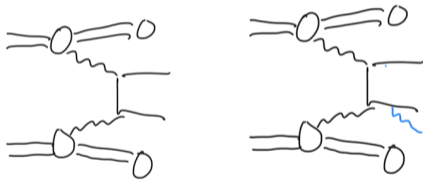
\Rightarrow Can apply photoproduction framework with all these beams!



$\gamma\gamma \rightarrow \mu^+\mu^-$ in proton-proton collisions

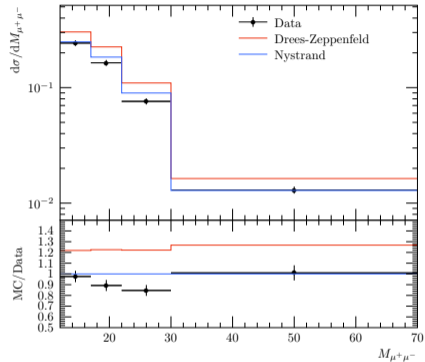
Elastic-elastic contribution

- Photons have small k_T proportional to Q^2
- Muons almost back-to-back ($A_{\text{co}} \approx 0$)
- Small effect from FSR



Clean process to calibrate flux

- Reasonable agreement with ATLAS data using EPA, DZ can be improved

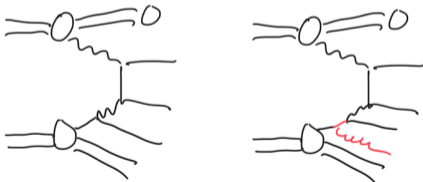


[ATLAS: PLB 777 (2018) 303-323]

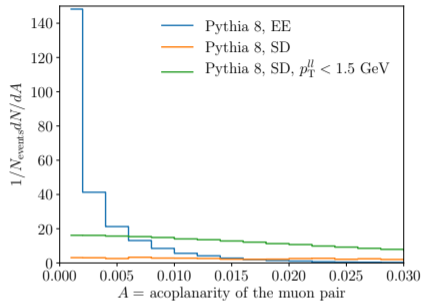
$\gamma\gamma \rightarrow \mu^+\mu^-$ in proton-proton collisions

Single-dissociative contribution

- Other γ from elastic flux, other as a part of DGLAP evolved proton PDFs
- Dissociative side will get primordial- k_T sampled from gaussian with width $\mathcal{O}(\text{GeV})$
- Also QCD ISR generated, significant p_T



- Cuts on $p_T^{\mu\mu}$ suppress events with ISR

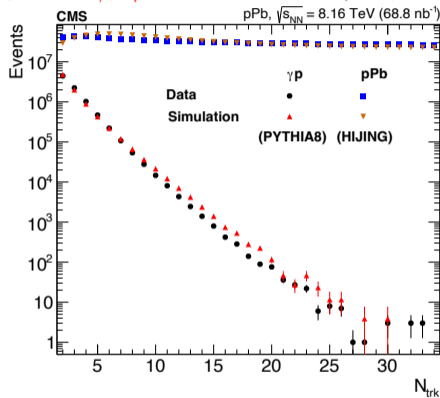


Double dissociative

- Both photons from PDFs with primordial- k_T and ISR
- ⇒ Large acoplanarity

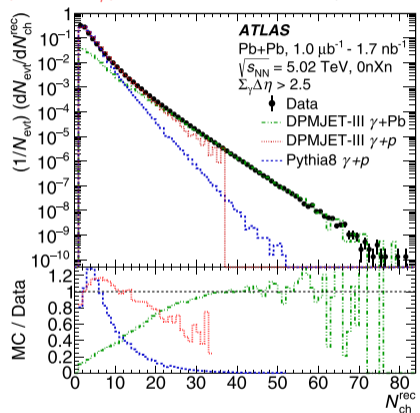
Ultra-peripheral heavy-ion collisions

$(\text{Pb} \rightarrow \gamma)+p$: [CMS: Murillo Quijada, QM2022]



- Multiplicities well reproduced with γp

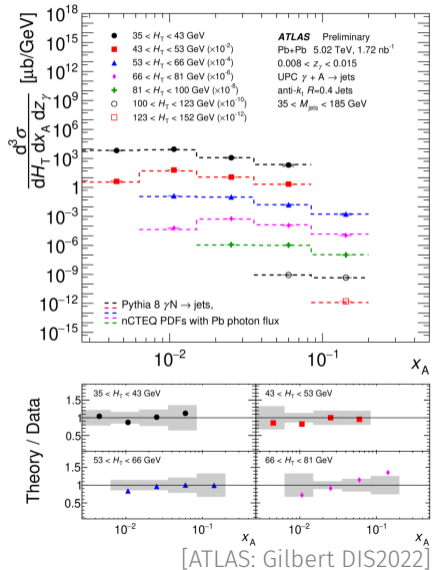
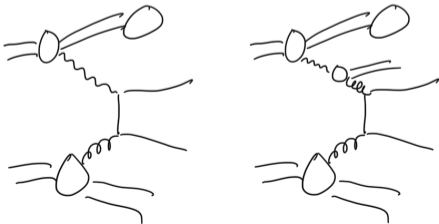
$(\text{Pb} \rightarrow \gamma)+\text{Pb}$: [ATLAS: PRC 104, 014903 (2021)]



- High multiplicities missed with γp
⇒ Multi-nucleon interactions

Dijets in ultra-peripheral heavy-ion collisions

- Novel constraints for nuclear PDFs, x_A to estimate probed nuclear x
- Pythia setup with nucleon target only
 \Rightarrow Not a realistic background for jet reconstruction
- Good agreement out of the box when accounting both direct and resolved



Hard diffraction in photoproduction

- Process with a hard scale, described with a colour-neutral Pomeron (IP) exchange
- Experimentally identified from rapidity gap

Factorization of the diffractive cross section

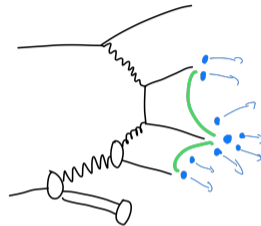
- Direct: Pomeron flux and diffractive PDFs (dPDFs)

$$d\sigma_{\text{direct}}^{2\text{jets}} = f_{\gamma}^b(x) \otimes d\sigma^{\gamma j \rightarrow 2\text{jets}} \otimes f_j^{\text{IP}}(z_{\text{IP}}, \mu^2) \otimes f_{\text{IP}}^{\text{D}}(x_{\text{IP}}, t)$$

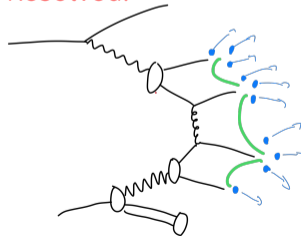
- Resolved: photon PDFs

$$d\sigma_{\text{resolved}}^{2\text{jets}} = f_{\gamma}^b(x) \otimes f_i^{\gamma}(x_{\gamma}, \mu^2) \otimes d\sigma^{ij \rightarrow 2\text{jets}} \otimes f_j^{\text{IP}}(z_{\text{IP}}, \mu^2) \otimes f_{\text{IP}}^{\text{D}}(x_{\text{IP}}, t)$$

Direct:



Resolved:



Hard diffraction in PYTHIA 8

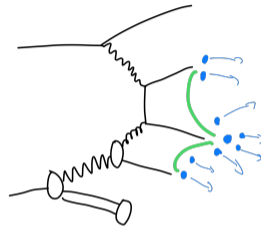
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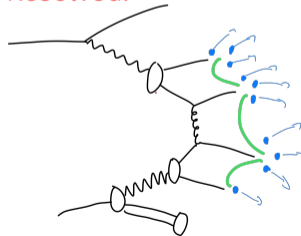
Dynamical rapidity gap survival model

1. Generate diffractive events with dPDFs (PDF)

Direct:



Resolved:



Hard diffraction in PYTHIA 8

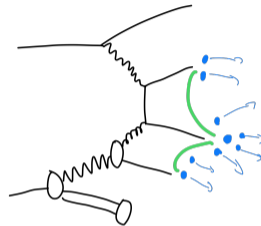
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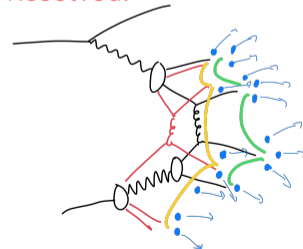
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2. Reject events where MPIs in γp system (MPI)

Direct:



Resolved:



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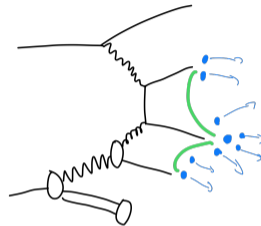
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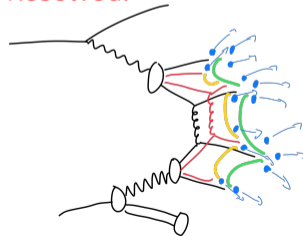
Dynamical rapidity gap survival model

1. Generate diffractive events with dPDFs (PDF)
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3. Evolve γIP system, allow MPIs

Direct:



Resolved:



Hard diffraction in PYTHIA 8

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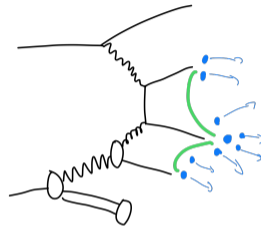
Implemented from PYTHIA 8.235 onwards

[I.H. and C.O. Rasmussen, EPJC 79 (2019) no.5, 413]

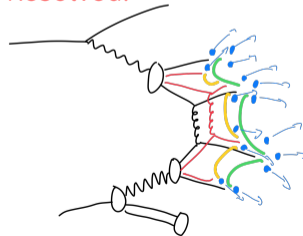
Same idea applied for pp collisions at the LHC

[C.O. Rasmussen and T. Sjöstrand, JHEP 1602 (2016) 142]

Direct:

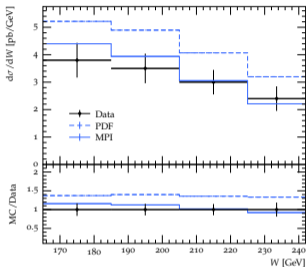


Resolved:

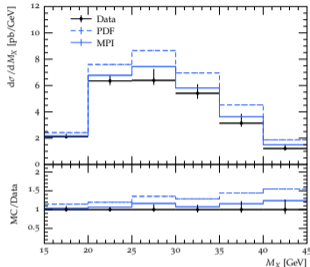


Comparisons to HERA data

H1: [EPJC 51 (2007) 549]



ZEUS: [EPJC 55 (2008) 177]



- PDF selection overshoots the data by 20–50 %
- Impact of the MPI rejection increases with W
- Stronger suppression in H1 analysis due to looser cuts on E_T^{jets} and $x_{\text{ip}} \Rightarrow$ More MPIs

Cuts

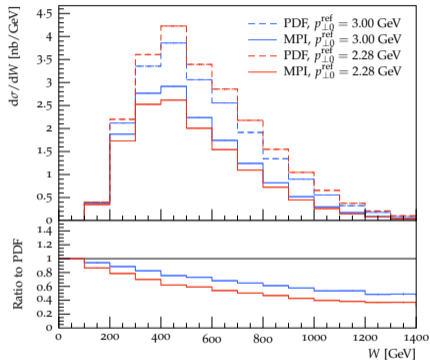
	H1	ZEUS
Q_{max}^2 [GeV ²]	0.01	1.0
$E_{T,\text{min}}^{\text{jet1}}$ [GeV]	5.0	7.5
$E_{T,\text{min}}^{\text{jet2}}$ [GeV]	4.0	6.5
$x_{\text{ip}}^{\text{max}}$	0.03	0.025

PYTHIA setup

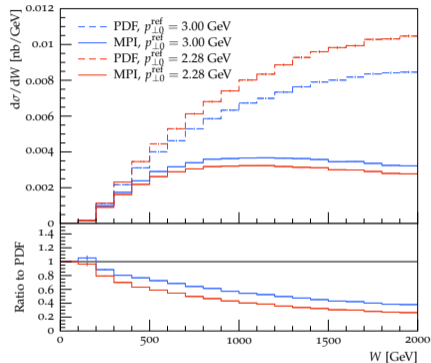
- dPDFs from H1 fit B LO
- γ PDFs from CJKL
- $p_{T0}^{\text{ref}} = 3.00$ GeV/c
(Tuned to inclusive charged particle data from γp at HERA)

Predictions for diffractive dijets in UPC

pPb $\sqrt{s} = 5.0$ TeV



pp $\sqrt{s} = 13$ TeV



- Extended W range wrt. HERA, especially in pp (harder flux)
- Stronger suppression from MPIs than at HERA
- ⇒ Ideal process to study factorization-breaking effects in hard diffraction

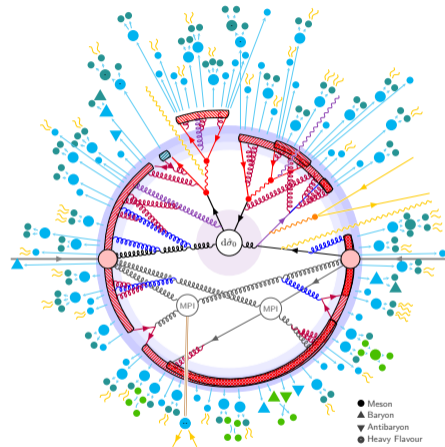
Summary & Outlook

Photon-induced processes in Pythia 8.3

- Photoproduction framework tested against HERA data
- Can be applied to purely hadronic collisions with appropriate fluxes
 - Fluxes in place for leptons, protons and heavy nuclei
 - Possibility to feed in externally provided flux

Outlook

- Subsequent resolved-photon nucleon interactions for $\gamma+A$ (Angantyr model)

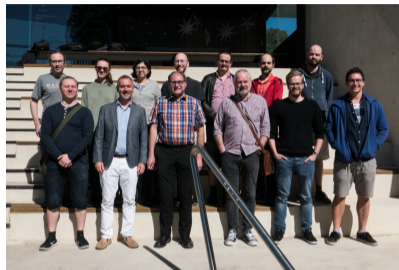


[figure by P. Skands]

Backup slides

PYTHIA Collaboration

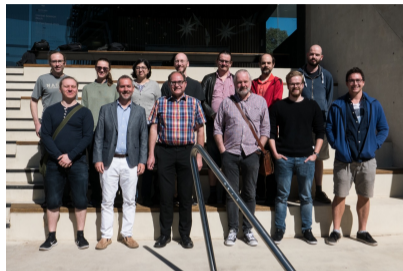
- Christian Bierlich (Lund University)
- Nishita Desai (TIFR, Mumbai)
- Leif Gellersen (Lund University)
- Ilkka Helenius (University of Jyväskylä)
- Philip Ilten (University of Cincinnati)
- Leif Lönnblad (Lund University)
- Stephen Mrenna (Fermilab)
- Stefan Prestel (Lund University)
- Christian Preuss (ETH Zurich)
- Torbjörn Sjöstrand (Lund University)
- Peter Skands (Monash University)
- Marius Utheim (University of Jyväskylä)
- Rob Verheyen (University College London)



[Pythia meeting in Monash 2019]

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[Pythia meeting in Monash 2019]

- Spokesperson
- Codemaster
- Webmaster

<https://pythia.org>
authors@pythia.org

Event generation in DIS with PYTHIA 8

Hard scattering

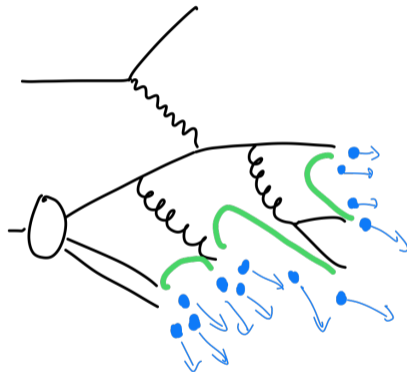
- Convolution between PDFs and matrix element (ME) for partonic scattering

Parton shower

- Final state radiation (FSR)
- Initial state radiation (ISR) for hadron
- QED emissions from leptons (omitted)

Hadronization

- String hadronization with colour reconnections
- Decays to stable hadrons



Alternative shower model `dipoleRecoil`

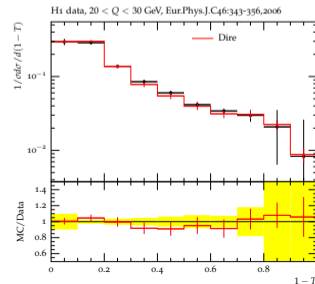
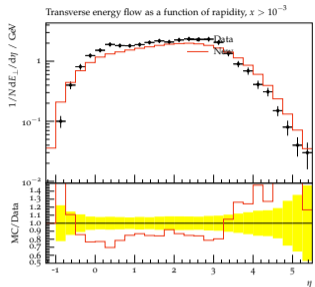
[B. Cabouat and T. Sjöstrand, EPJC 78 (2018 no.3, 226)]

- No PS recoil for the scattered lepton
- Reasonable description of single-particle properties, such as transverse energy flow
- Results based on tune with the default global-recoil shower

Completely new shower `DIRE`

[S. Höche, S. Prestel, EPJC 75 (2015) no.9, 461]

- Correct soft-gluon interference at lowest order
- Inclusive NLO corrections to collinear splittings
- Good agreement with HERA data e.g. for thrust T



Evolution equation and PDFs for resolved photons

DGLAP equation for photons

- Additional term due to $\gamma \rightarrow q\bar{q}$ splittings

$$\frac{\partial f_i^\gamma(x, Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{em}}{2\pi} e_i^2 P_{i\gamma}(x) + \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{dz}{z} P_{ij}(z) f_j(x/z, Q^2)$$

where $P_{i\gamma}(x) = 3(x^2 + (1-x)^2)$ for quarks, 0 for gluons (LO)

- Resulting PDFs has **point-like** (or anomalous) and **hadron-like** components

$$f_i^\gamma(x, Q^2) = f_i^{\gamma, pl}(x, Q^2) + f_i^{\gamma, had}(x, Q^2)$$

- $f_i^{\gamma, pl}$: Calculable from perturbative QCD
- $f_i^{\gamma, had}$: Requires non-perturbative input fixed in a global analysis

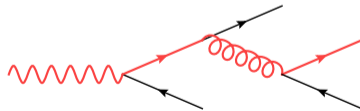
Evolution equation and ISR for resolved photons

ISR probability based on DGLAP evolution

- Add a term corresponding to $\gamma \rightarrow q\bar{q}$ to (conditional) ISR probability

$$d\mathcal{P}_{a\leftarrow b} = \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} \frac{x' f_a^\gamma(x', Q^2)}{x f_b^\gamma(x, Q^2)} P_{a\rightarrow bc}(z) dz + \frac{dQ^2}{Q^2} \frac{\alpha_{em}}{2\pi} \frac{e_b^2 P_{\gamma\rightarrow bc}(x)}{f_b^\gamma(x, Q^2)}$$

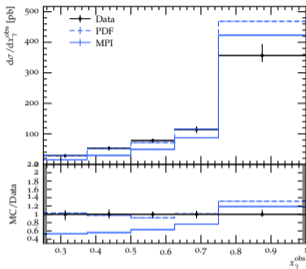
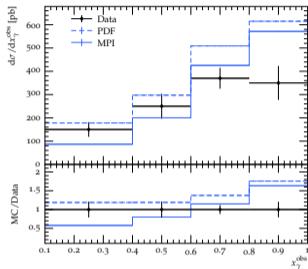
- Corresponds to ending up to the beam photon during evolution
 - \Rightarrow Parton originated from the point-like part of the PDFs
 - No further ISR or MPIs below the scale of the splitting
 - No need for beam remnants



Comparisons to HERA data

H1: [EPJC 51 (2007) 549]

ZEUS: [EPJC 55 (2008) 177]



- Stronger suppression at low- x_{γ}^{obs} (more MPIs)
- ZEUS cuts select events at high- x_{γ}^{obs} region
- Some theoretical uncertainty from γ PDFs, dPDFs and scale variation

Cuts

H1

ZEUS

Q_{max}^2 [GeV²]

0.01

1.0

$E_{T,\text{min}}^{\text{jet1}}$ [GeV]

5.0

7.5

$E_{T,\text{min}}^{\text{jet2}}$ [GeV]

4.0

6.5

$x_{\text{ip}}^{\text{max}}$

0.03

0.025

χ^2 analysis

PDF

MPI

H1

5.2

1.4

ZEUS

9.6

5.1

H1 & ZEUS

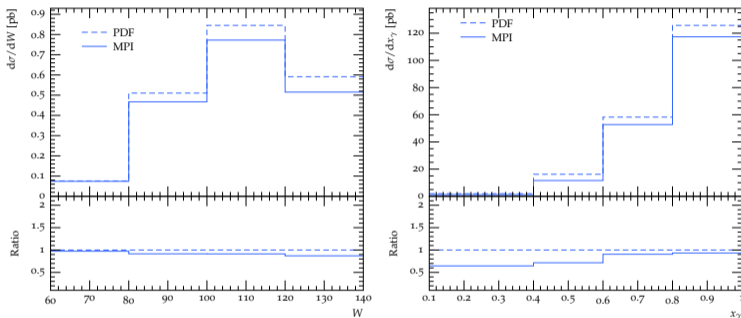
7.6

3.4

(with all data points)

Predictions for EIC

Repeat the H1 analysis at EIC kinematics ($E_e = 18$ GeV, $E_p = 275$ GeV)



- Only up to $\sim 10\%$ effects in the considered W range
 - Noticeable suppression only at low x_γ where cross section small
- \Rightarrow Available energy and kinematical cuts for diffraction push the kinematics to region where only little room for MPIs ($E_T^{\text{jet}1} > 5.0$ GeV, $E_T^{\text{jet}2} > 4.0$ GeV)

Intermediate Q^2 region

Solid theory for $Q^2 = 0$ and at high Q^2

⇒ What happens in between?

Pythia 6 (inspired) model (\neq Pythia 8)

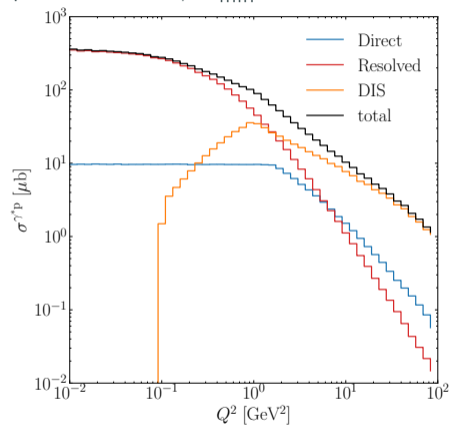
- Select suitable scales and suppress contributions by hand

$$\sigma_{\text{tot}}^{\gamma^*p} = \tilde{\sigma}_{\text{DIS}}^{\gamma^*p} \exp\left[-\frac{\tilde{\sigma}_{\text{Dir}}^{\gamma^*p}}{\tilde{\sigma}_{\text{DIS}}^{\gamma^*p}}\right] + \tilde{\sigma}_{\text{Dir}}^{\gamma^*p} + \tilde{\sigma}_{\text{Res}}^{\gamma^*p}$$

where

- $\tilde{\sigma}_{\text{DIS}}^{\gamma^*p} = \left[\frac{Q^2}{Q^2+m_\rho^2}\right]^2 \sigma_{\text{DIS}}^{\gamma^*p}$
- $\tilde{\sigma}_{\text{Res}}^{\gamma^*p} = \sigma_{\text{Res}}^{\gamma^*p} \left[\frac{m_\rho^2}{m_\rho^2+Q^2}\right]^2 \left[\frac{W^2}{W^2+Q^2}\right]^n$
- $\tilde{\sigma}_{\text{Dir}}^{\gamma^*p} = \sigma_{\text{Dir}}^{\gamma^*p} (\hat{p}_{\text{T,min}} = \max(Q, p_{\text{T,min}}))$
 $p_{\text{T,min}} = 1.3 \text{ GeV}, n = 3, m_\rho = 0.7755 \text{ GeV}$

$\sqrt{s} = 318 \text{ GeV}, W_{\text{min}} = 100 \text{ GeV}$:



Intermediate: $0.5 \lesssim Q^2 \lesssim 5.0 \text{ GeV}^2$

Intermediate Q^2 region

Solid theory for $Q^2 = 0$ and at high Q^2

⇒ What happens in between?

Pythia 6 (inspired) model (\neq Pythia 8)

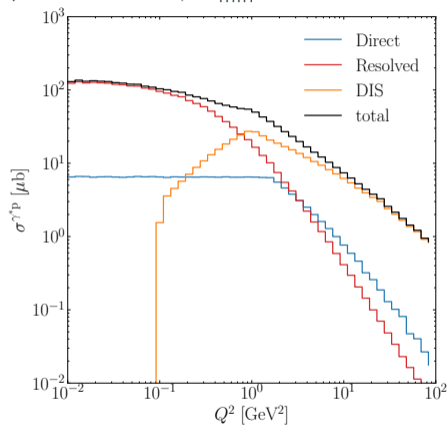
- Select suitable scales and suppress contributions by hand

$$\sigma_{\text{tot}}^{\gamma^*p} = \tilde{\sigma}_{\text{DIS}}^{\gamma^*p} \exp \left[-\frac{\tilde{\sigma}_{\text{Dir}}^{\gamma^*p}}{\tilde{\sigma}_{\text{DIS}}^{\gamma^*p}} \right] + \tilde{\sigma}_{\text{Dir}}^{\gamma^*p} + \tilde{\sigma}_{\text{Res}}^{\gamma^*p}$$

where

- $\tilde{\sigma}_{\text{DIS}}^{\gamma^*p} = \left[\frac{Q^2}{Q^2 + m_\rho^2} \right]^2 \sigma_{\text{DIS}}^{\gamma^*p}$
- $\tilde{\sigma}_{\text{Res}}^{\gamma^*p} = \sigma_{\text{Res}}^{\gamma^*p} \left[\frac{m_\rho^2}{m_\rho^2 + Q^2} \right]^2 \left[\frac{W^2}{W^2 + Q^2} \right]^n$
- $\tilde{\sigma}_{\text{Dir}}^{\gamma^*p} = \sigma_{\text{Dir}}^{\gamma^*p} (\hat{p}_{\text{T,min}} = \max(Q, p_{\text{T,min}}))$
 $p_{\text{T,min}} = 1.3 \text{ GeV}, n = 3, m_\rho = 0.7755 \text{ GeV}$

$\sqrt{s} = 140 \text{ GeV}, W_{\text{min}} = 10 \text{ GeV}$:



Intermediate: $0.3 \lesssim Q^2 \lesssim 3.0 \text{ GeV}^2$

Intermediate Q^2 region

Solid theory for $Q^2 = 0$ and at high Q^2

⇒ What happens in between?

Pythia 6 (inspired) model (\neq Pythia 8)

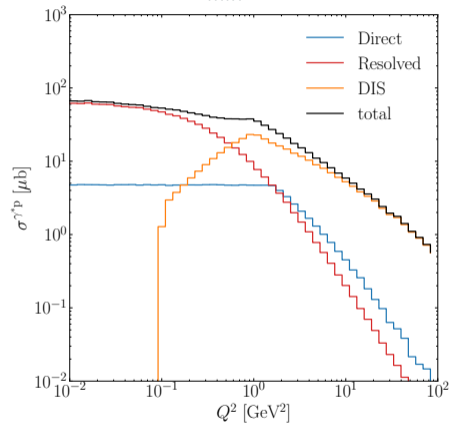
- Select suitable scales and suppress contributions by hand

$$\sigma_{\text{tot}}^{\gamma^*p} = \tilde{\sigma}_{\text{DIS}}^{\gamma^*p} \exp \left[-\frac{\tilde{\sigma}_{\text{Dir}}^{\gamma^*p}}{\tilde{\sigma}_{\text{DIS}}^{\gamma^*p}} \right] + \tilde{\sigma}_{\text{Dir}}^{\gamma^*p} + \tilde{\sigma}_{\text{Res}}^{\gamma^*p}$$

where

- $\tilde{\sigma}_{\text{DIS}}^{\gamma^*p} = \left[\frac{Q^2}{Q^2 + m_\rho^2} \right]^2 \sigma_{\text{DIS}}^{\gamma^*p}$
- $\tilde{\sigma}_{\text{Res}}^{\gamma^*p} = \sigma_{\text{Res}}^{\gamma^*p} \left[\frac{m_\rho^2}{m_\rho^2 + Q^2} \right]^2 \left[\frac{W^2}{W^2 + Q^2} \right]^n$
- $\tilde{\sigma}_{\text{Dir}}^{\gamma^*p} = \sigma_{\text{Dir}}^{\gamma^*p} (\hat{p}_{T,\text{min}} = \max(Q, p_{T,\text{min}}))$
 $p_{T,\text{min}} = 1.3 \text{ GeV}, n = 3, m_\rho = 0.7755 \text{ GeV}$

$\sqrt{s} = 85 \text{ GeV}, W_{\text{min}} = 10 \text{ GeV}$:



Intermediate: $0.2 \lesssim Q^2 \lesssim 2.0 \text{ GeV}^2$